

Collisional Excitation of PO^+ by para- H_2



Towards the Accurate Modeling of Phosphorus Abundance in the ISM

PCMI2024-Bordeaux
Colloque du Programme Physique Chimie du Milieu Interstellaire

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COLLEXISM



PO⁺ and H₂ system

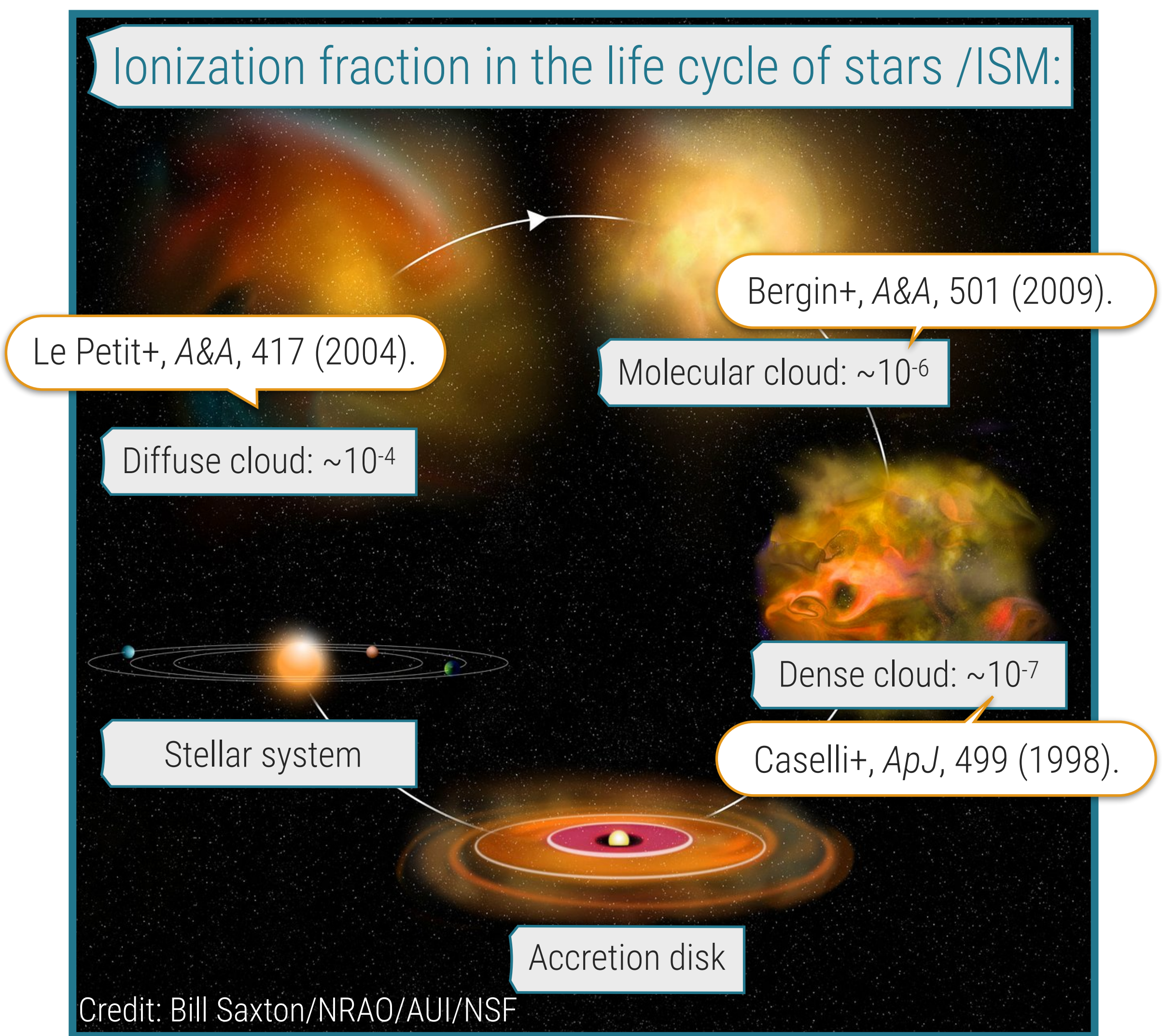
Why Ions?

The ionization degree is a **diagnostic** tool of the **evolutionary stage** of astrophysical objects.

Ions play a prominent role in **chemical reactivity**.

$$x(e^-) = n(e^-)/n(H_2)$$

Ionization fraction in the life cycle of stars /ISM:



PO⁺ and H₂ system

Why Phosphorus?

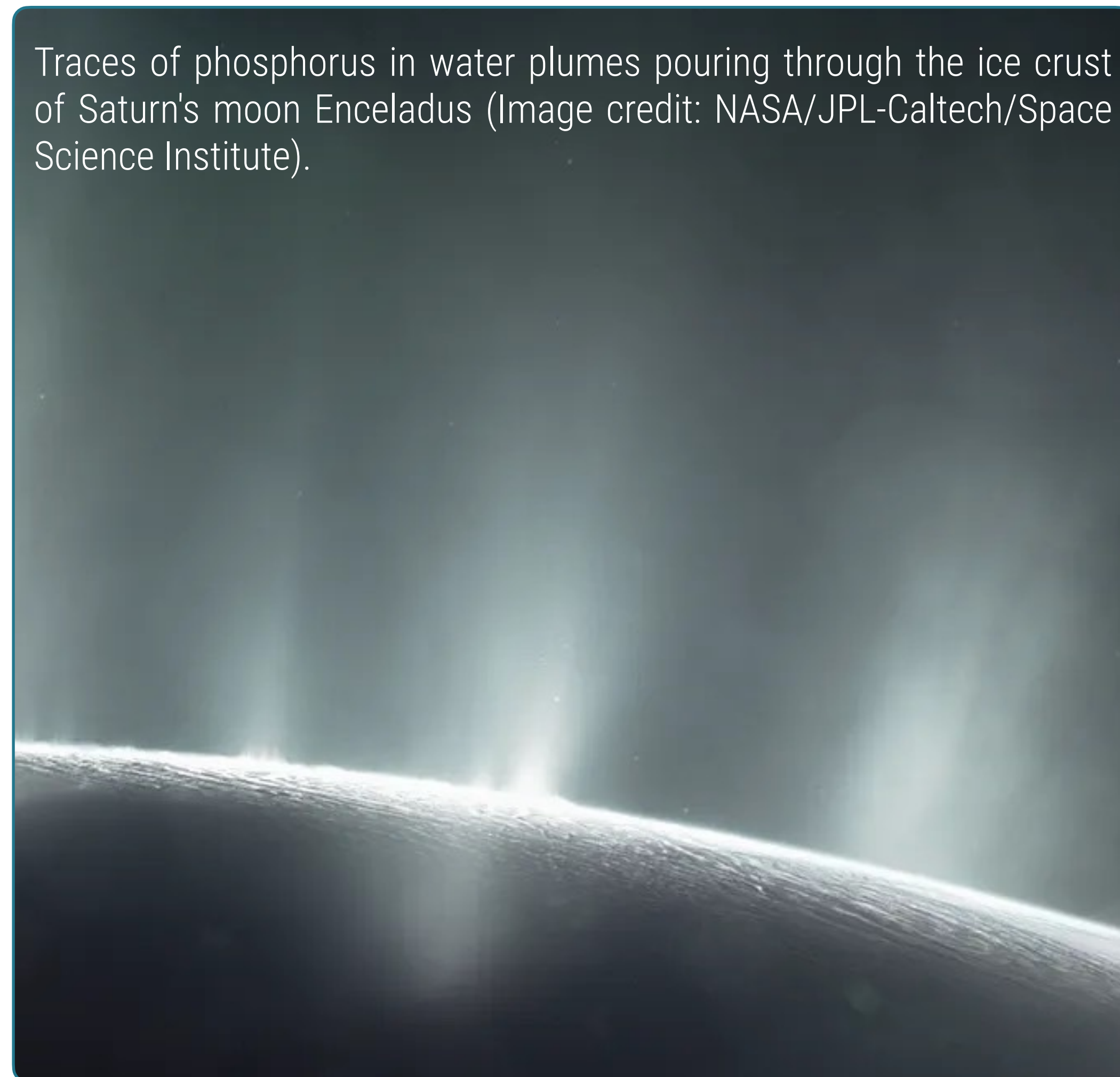
B. K. Pearce+, *Proc. Natl. Acad. Sci*, **114**, (2017).

Element of pivotal importance from a **prebiotic** point of view. It is ubiquitous in our planet and P-bearing compounds have been found also in many environments outside Earth.

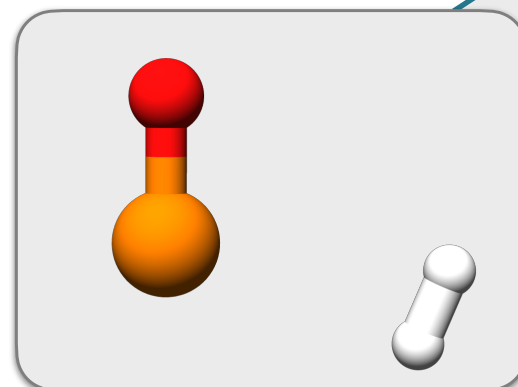
However, only **few** a P-bearing molecules have been **identified in the ISM**:

- **Depletion** onto interstellar grains;
- **Paucity of observational data** and chemical modeling studies.

Traces of phosphorus in water plumes pouring through the ice crust of Saturn's moon Enceladus (Image credit: NASA/JPL-Caltech/Space Science Institute).



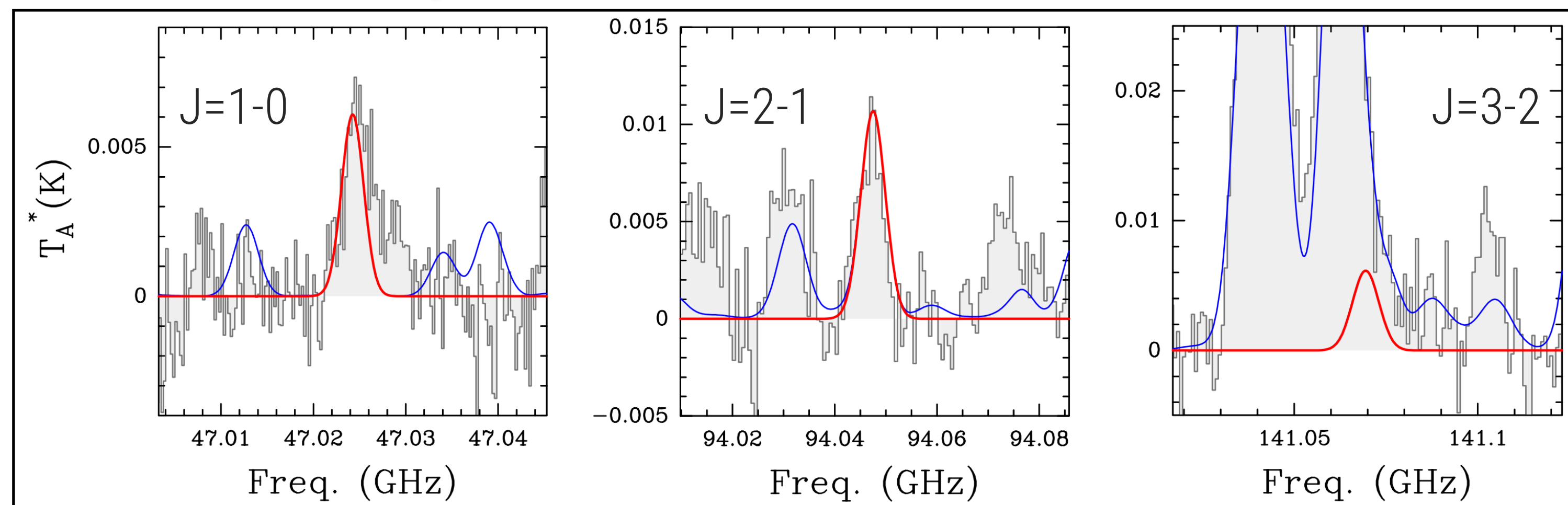
PO⁺ and H₂ system



Target species: PO⁺ recently detected in the molecular cloud G+0.693-0.02 where, due to the low density of H₂, LTE condition is not achieved.

Collider: in ISM, the most abundant perturber is H₂.

V. M. Rivilla, et al. (2022), *FSPAS*, **9**, 829288.



G+0.693-0.027 molecular cloud:

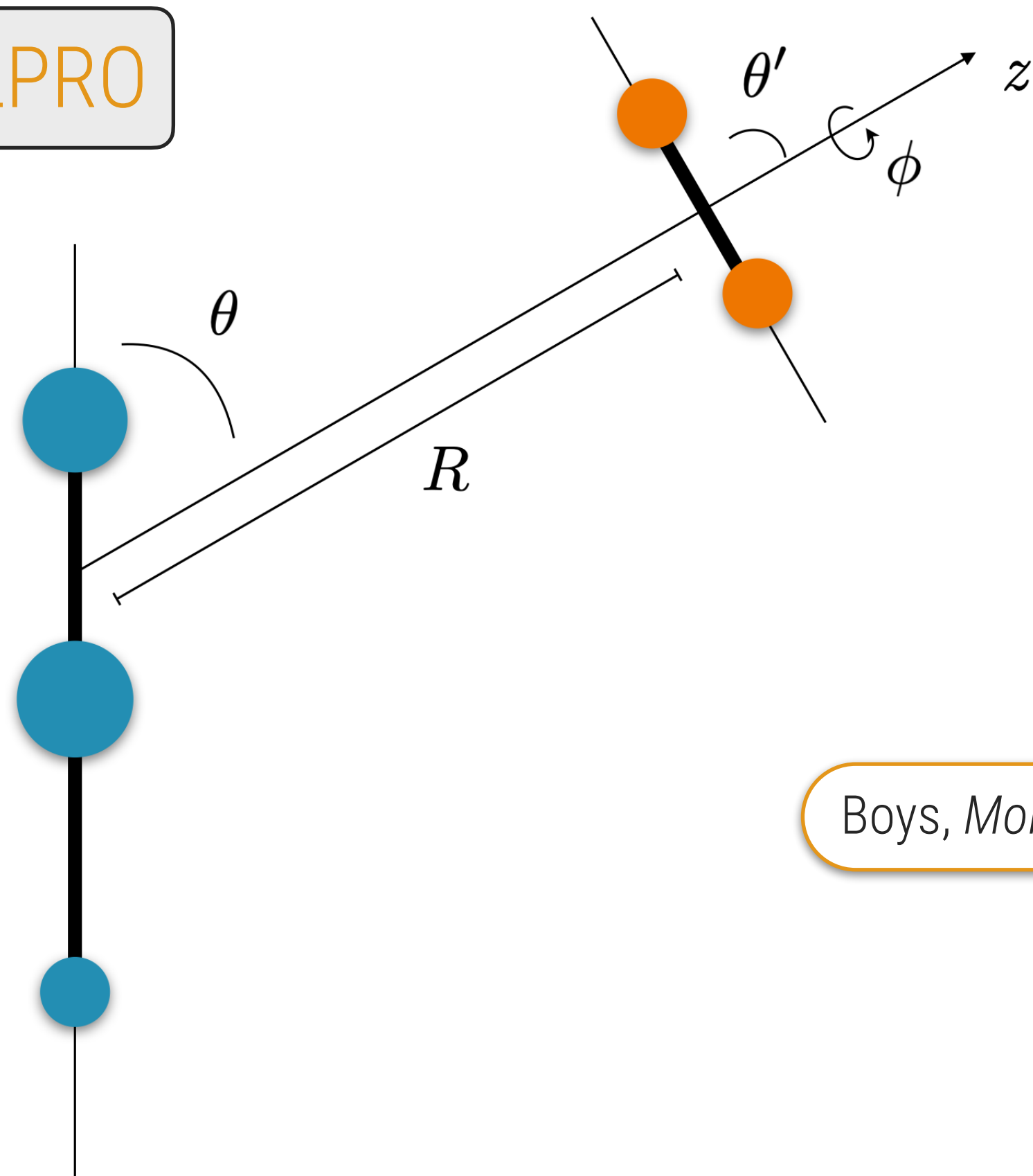
- Very low densities of H₂ ($\sim 1 \times 10^4 \text{ cm}^{-3}$);
- High kinetic temperatures ($\sim 150 \text{ K}$);

S. Zeng, et al. (2018), *MNRAS*, **478**, 2962;
S. Zeng, et al. (2020), *MNRAS*, **497**, 4896.

Intermolecular PES



MOLPRO



Boys, *Mol.Phys.*, **19** (1970).

Single point energies:

CCSD(T)-F12/aug-cc-pV(Q+d)Z level of theory.

Geometries:

- The two colliders are considered as **rigid bodies** and only the interaction between them is taken into account.
- The body-fixed **Jacobi coordinates** frame is employed.

Interaction energies:

- *Ab initio* calculations over an irregular grid in the $\{R, \theta, \theta', \phi\}$ coordinates.
- Each energy has been corrected for the basis set superposition error (**BSSE**).

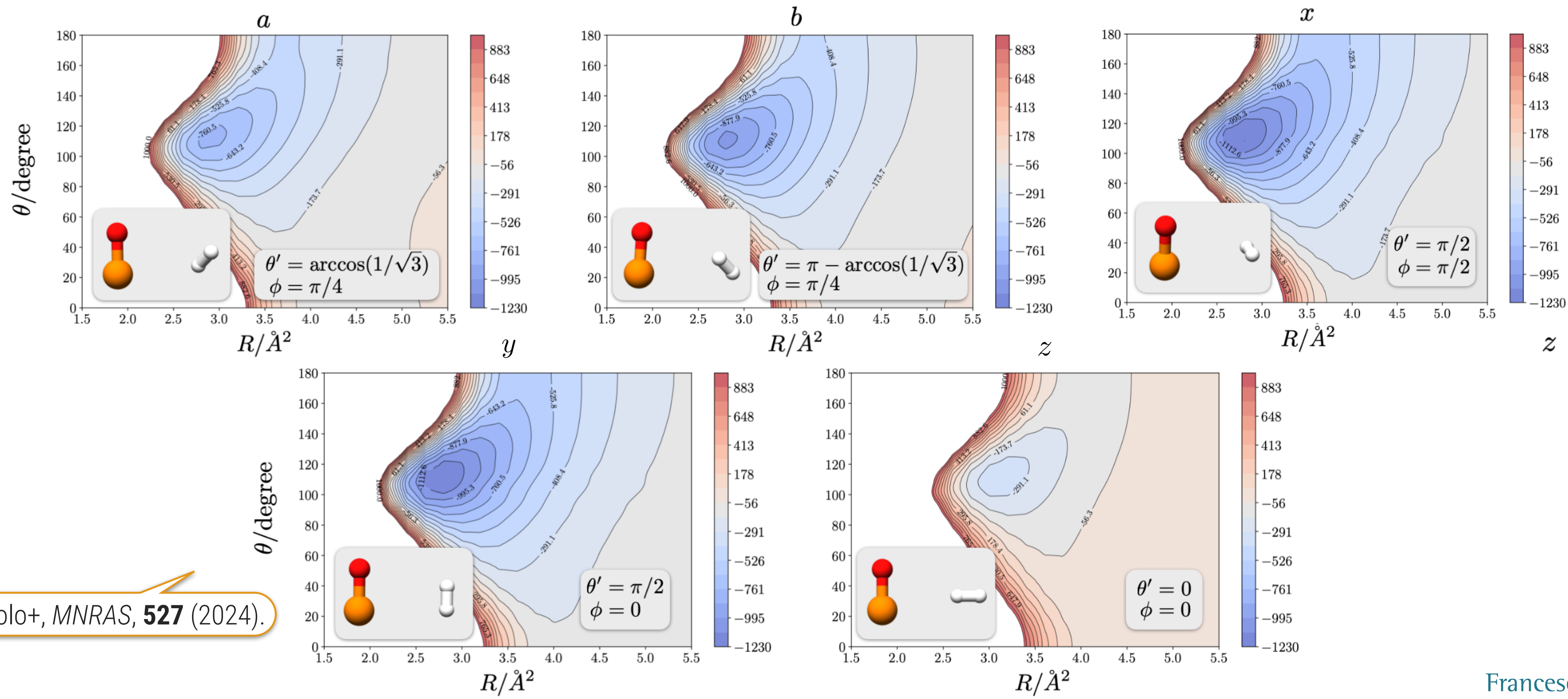
Fitting of the potential:

$$V(R, \theta, \theta', \phi) = \sum_{l_1 l_2 \mu} v_{l_1 l_2 \mu}(R) s_{l_1 l_2 \mu}(\theta, \theta', \phi)$$

Jacobi internal coordinates of a collisional system.

Fitting the 4D potential

$$V(R, \theta, \theta', \phi) = \sum_{l_1 l_2 \mu} v_{l_1 l_2 \mu}(R) s_{l_1 l_2 \mu}(\theta, \theta', \phi)$$



Tonolo+, MNRAS, **527** (2024).

Fitting the 4D potential

$$V(R, \theta, \theta', \phi) = \sum_{l_1 l_2 \mu} v_{l_1 l_2 \mu}(R) s_{l_1 l_2 \mu}(\theta, \theta', \phi)$$

where: $s_{l_1, l_2, \mu}(\theta, \theta', \phi) = A_{l_1} \left[\langle l_1 0 l_2 0 | l_1 l_2 \mu 0 \rangle P_{l_1 0}(\theta) P_{l_2 0}(\theta') + \sum_m (-1)^m 2 \langle l_1 m l_2 -m | l_1 l_2 \mu 0 \rangle P_{l_1 m}(\theta) P_{l_2 m}(\theta') \cos(m\phi) \right],$
 $P_{lm}(\theta) = Y_{lm}(\theta, \phi) \exp(-im\phi).$

Truncation of the expansion at the $l_2 < 2$ terms of the angular momentum of the collider:

The only basis function describing the dependence of the potential on the relative orientations of the collider are $Y_{00}, Y_{20}, Y_{21}, Y_{22}$.



Reliable choice of **five sets** of $\{\theta', \phi\}$ coordinates to describe the system:

$$V(R, \theta, \theta', \phi) = V_{iso} + \frac{1}{2} [V_{eff}(z) - V_{iso}] (3 \cos^2 \theta') + \frac{3}{2} [V_{eff}(a) - V_{eff}(b)] \sin \theta' \cos \theta' \cos \phi + \frac{1}{2} [V_{eff}(x) - V_{eff}(y)] \sin^2 \theta' \cos(2\phi)$$

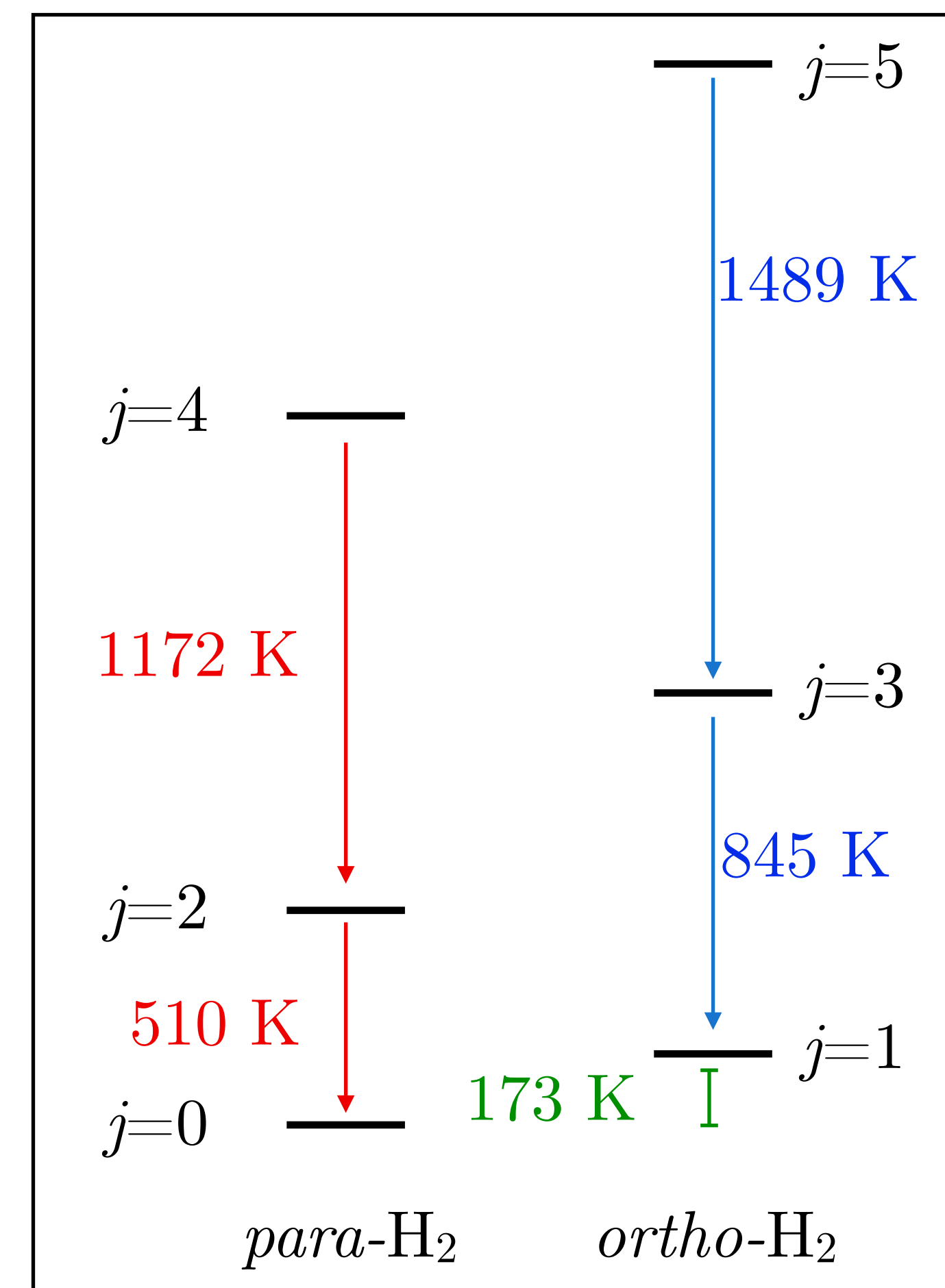
Spherical average approximation

If the **inclusion of the rotational structure** ($j \neq 0$) of the collider has a negligible effect on the cross sections' values, a spherical average of the potential with respect to five orientations of the collider can be performed:

$$V_{\text{av}}(R, \theta) = \sum_{l_1} v_{l_1}(R) P_{l_1}(\cos(\theta)),$$

$$V_{\text{av}}(R, \theta) = \frac{1}{7} [2(V(R, \theta, a) + V(R, \theta, b)) + (V(R, \theta, x) + V(R, \theta, y) + V(R, \theta, z))].$$

The collider is considered as a **structureless species**, behaving as a rotating sphere.



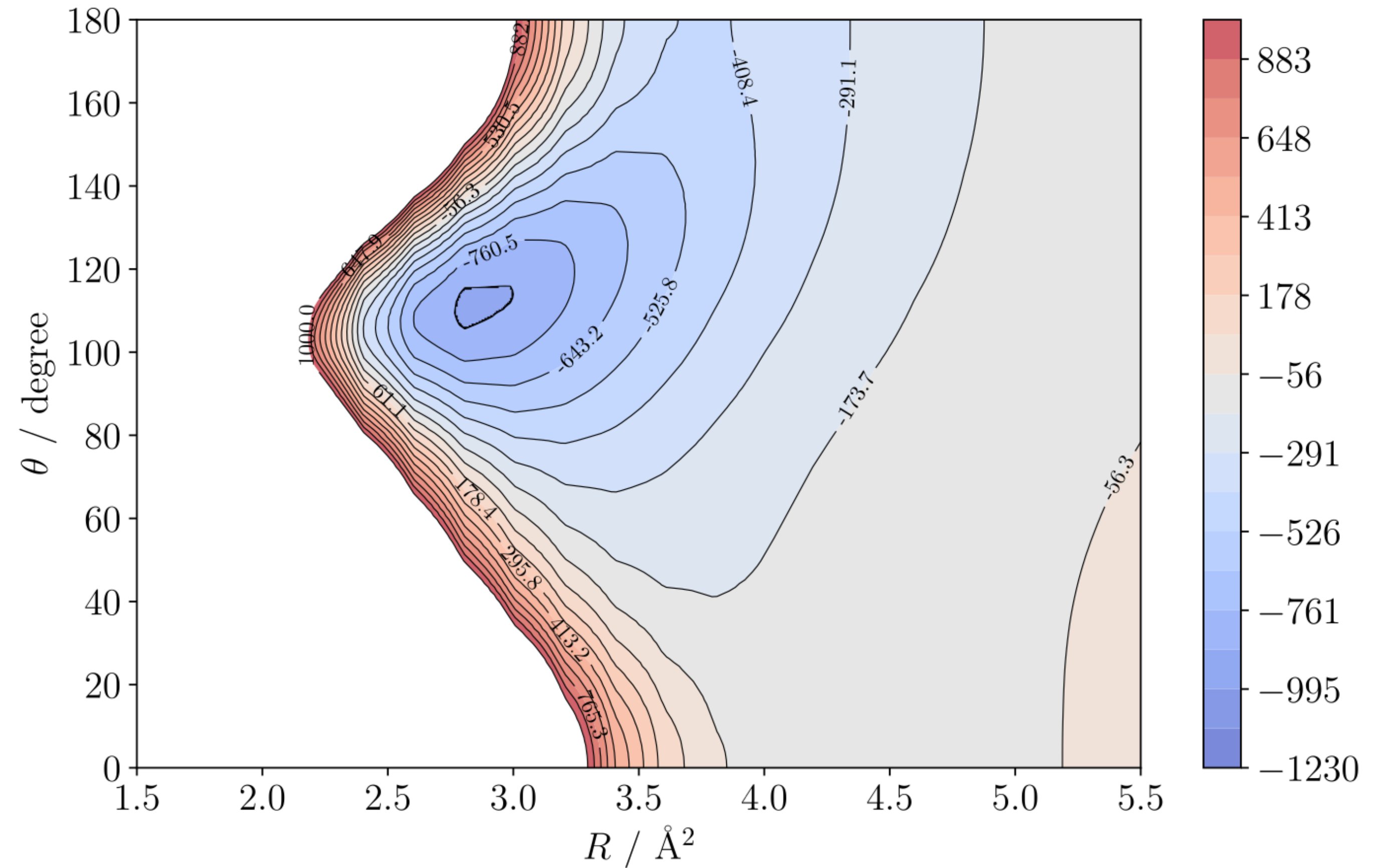
Lique+, A&A, **478** (2008);
Dumouchel+, JCP, **137** (2012).

Spherical average approximation

Tonolo+, *MNRAS*, **527** (2024).

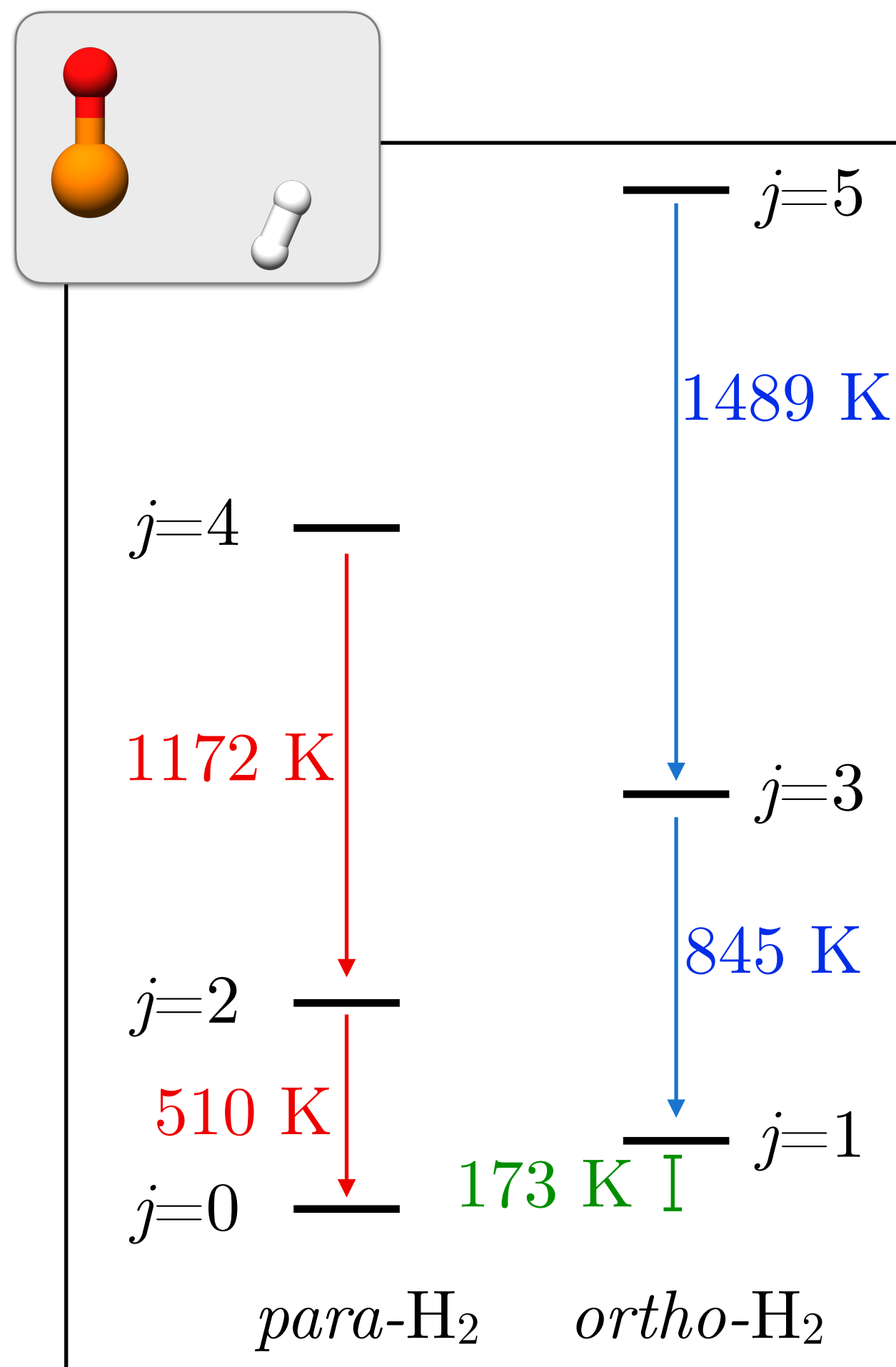
$j' \rightarrow j$	Cross sections / \AA^2		% Deviation 2D/4D
	2D	4D	
1 → 0	33.10	30.46	-8.68
2 → 0	17.43	19.16	9.02
3 → 0	9.32	8.48	-9.88
4 → 0	12.64	11.89	-6.29
5 → 0	11.43	11.68	2.10
6 → 0	10.55	11.76	10.30
2 → 1	26.61	24.55	-8.42
3 → 1	16.84	15.16	-11.08
4 → 1	11.75	12.82	8.38
5 → 1	11.57	11.19	-3.42
6 → 1	13.28	14.05	5.45
3 → 2	26.06	30.57	14.73
4 → 2	15.67	14.23	-10.14
5 → 2	12.61	12.92	2.39
6 → 2	12.77	13.55	5.76
4 → 3	27.30	28.84	5.32
5 → 3	14.64	14.90	1.77
6 → 3	13.89	12.39	-12.11
5 → 4	27.51	26.88	-2.34
6 → 5	15.64	14.98	-4.44
Average absolute % deviation			7.10

PO⁺/H₂ cross sections at $E_{\text{kin}}=150 \text{ cm}^{-1}$



Contour plot of the averaged potential over five $\{\theta, \phi\}$ orientations.

ortho- and para- states of the collider

 Tonolo+, *MNRAS*, **527** (2024).


In certain circumstances, collisions with the $j=1$ state of the collider can also have an impact on scattering calculations.

This is the case of the G+0.693-0.027 molecular cloud where PO⁺ was detected, with kinetic temperatures between 140 and 200 K.

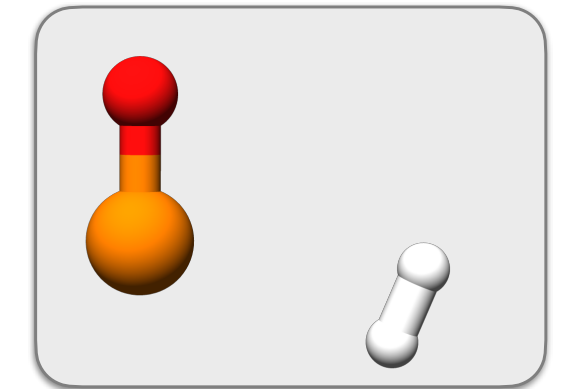
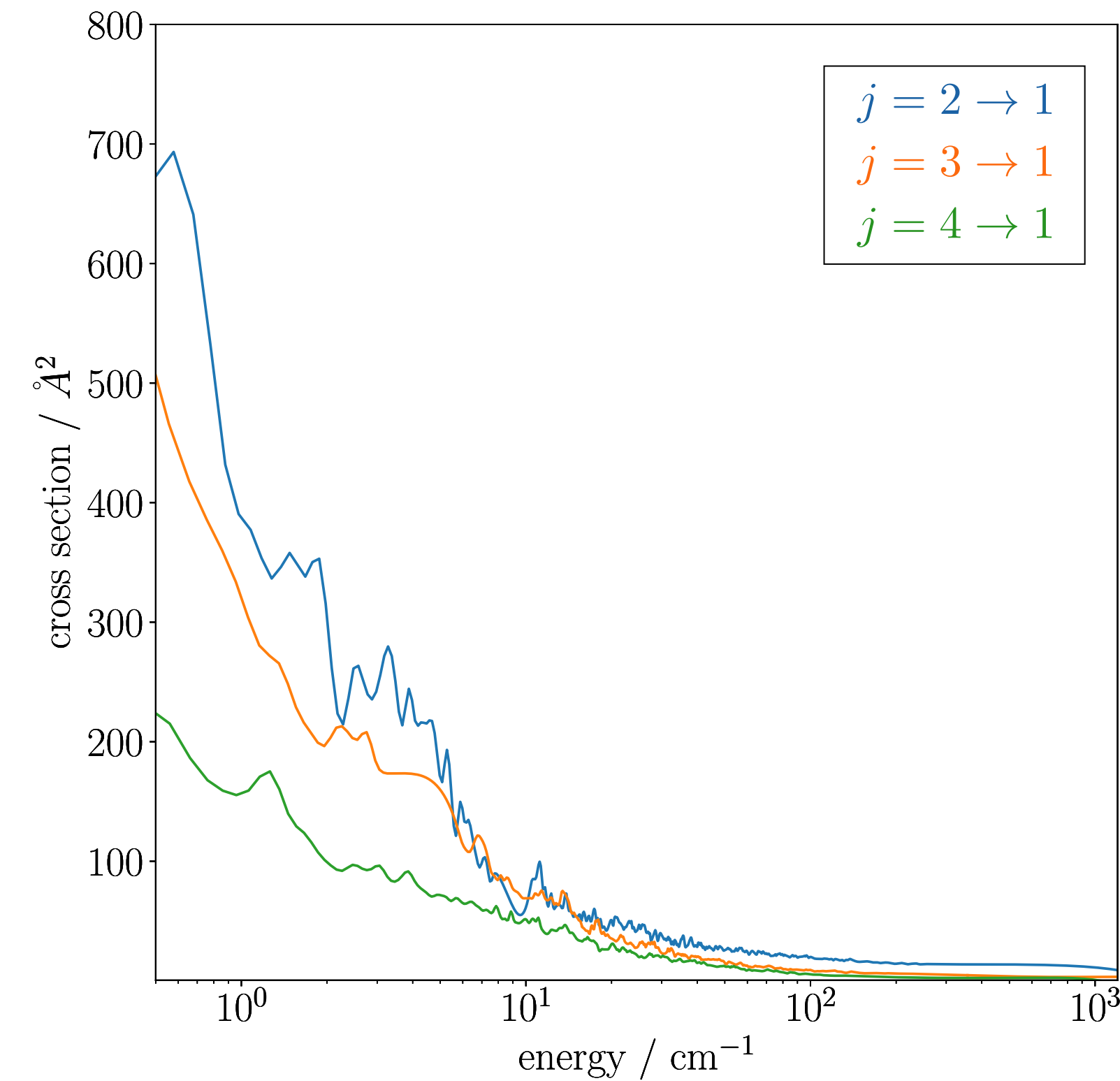
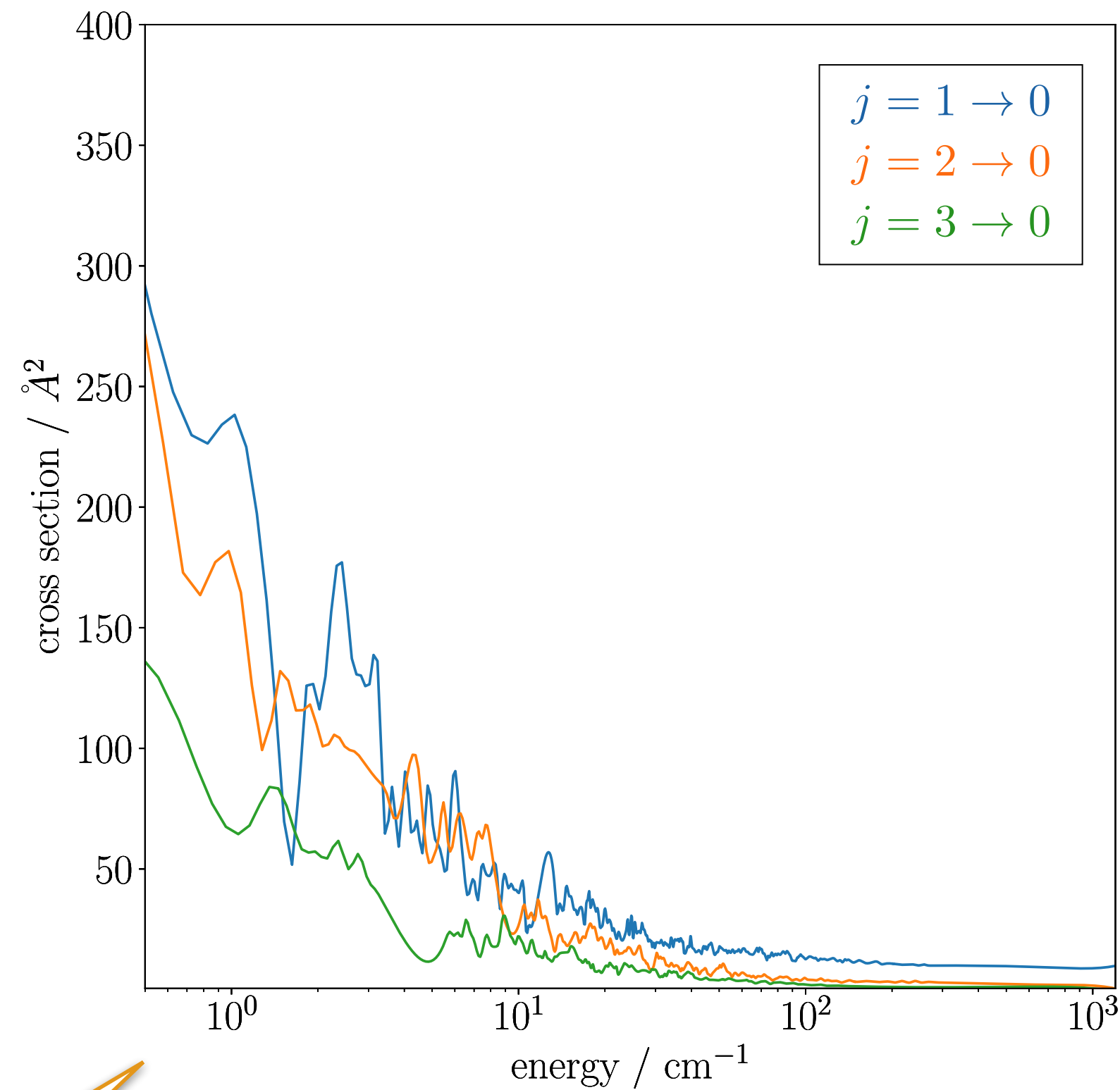
S. Zeng, et al. (2018), *MNRAS*, **478**, 2962;
S. Zeng, et al. (2020), *MNRAS*, **497**, 4896.

$j' \rightarrow j$	Cross sections / Å ²		% Deviation <i>o-/p-</i>
	$j=0$	$j=1$	
1 → 0	33.10	33.84	-2.19
2 → 0	17.43	20.46	-14.81
3 → 0	9.32	10.88	-14.35
4 → 0	12.64	14.39	-12.17
5 → 0	11.43	12.56	-9.02
6 → 0	10.55	12.26	-13.95
2 → 1	26.61	29.11	-8.57
3 → 1	16.84	19.70	-14.52
4 → 1	11.75	13.16	-10.72
5 → 1	11.57	13.92	-16.87
6 → 1	13.28	17.37	-23.56
3 → 2	26.06	27.96	-6.79
4 → 2	15.67	18.19	-13.87
5 → 2	12.61	14.26	-11.54
6 → 2	12.77	17.02	-24.93
4 → 3	27.30	31.06	-12.10
5 → 3	14.64	16.36	-10.54
6 → 3	13.89	16.93	-17.98
5 → 4	27.51	29.77	-7.59
6 → 5	15.64	17.08	-8.44
Average absolute % deviation			12.73

 $E_{\text{kin}} = 150 \text{ cm}^{-1}$

close coupling calculations

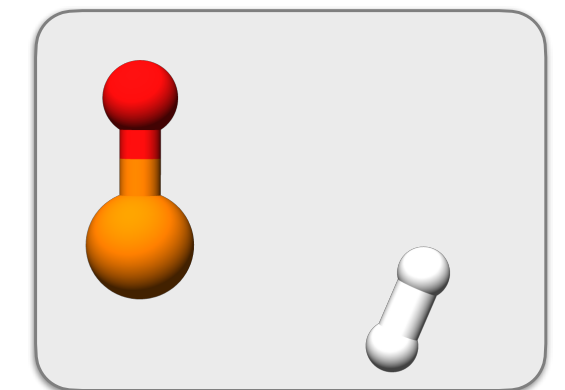
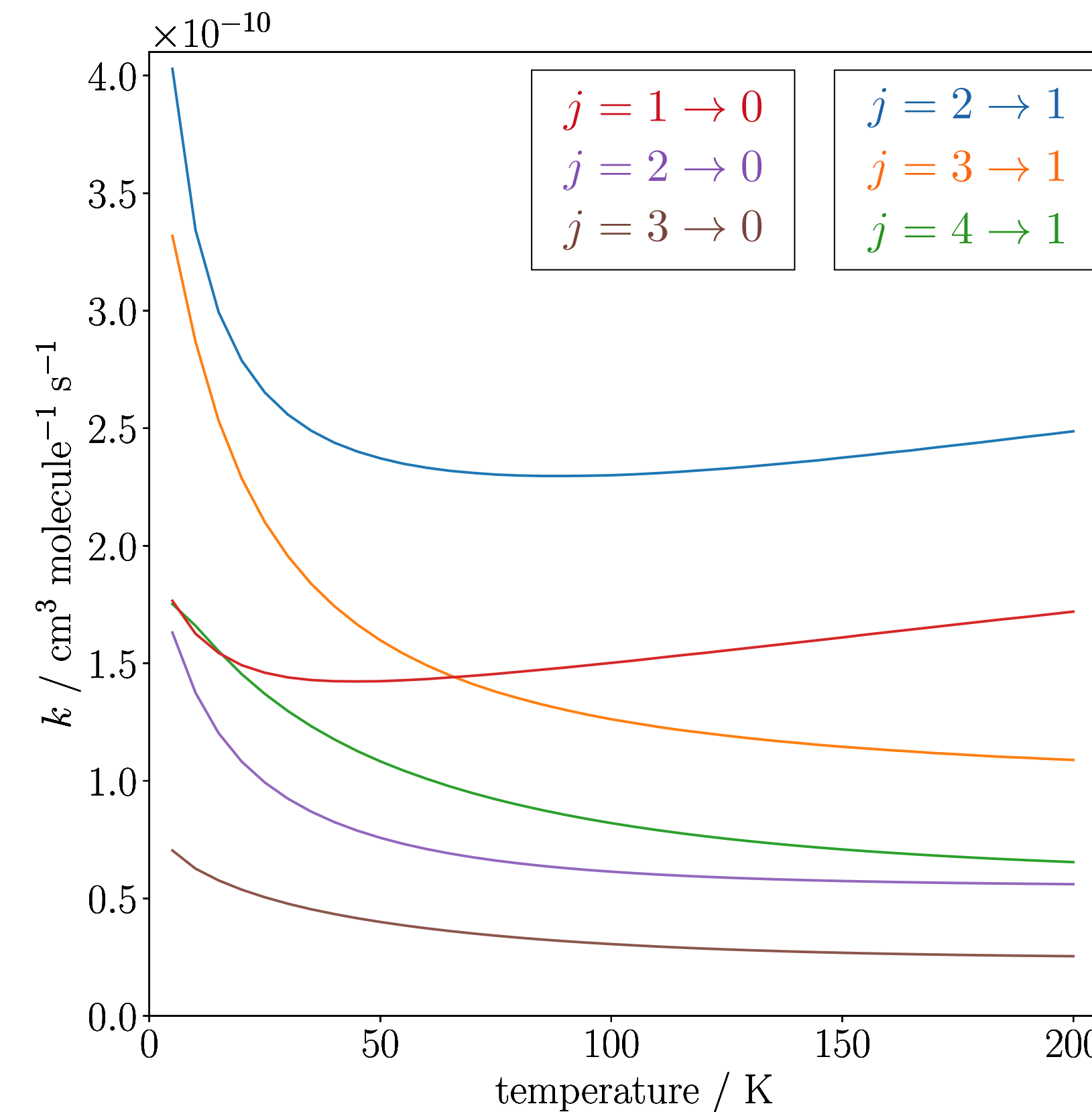
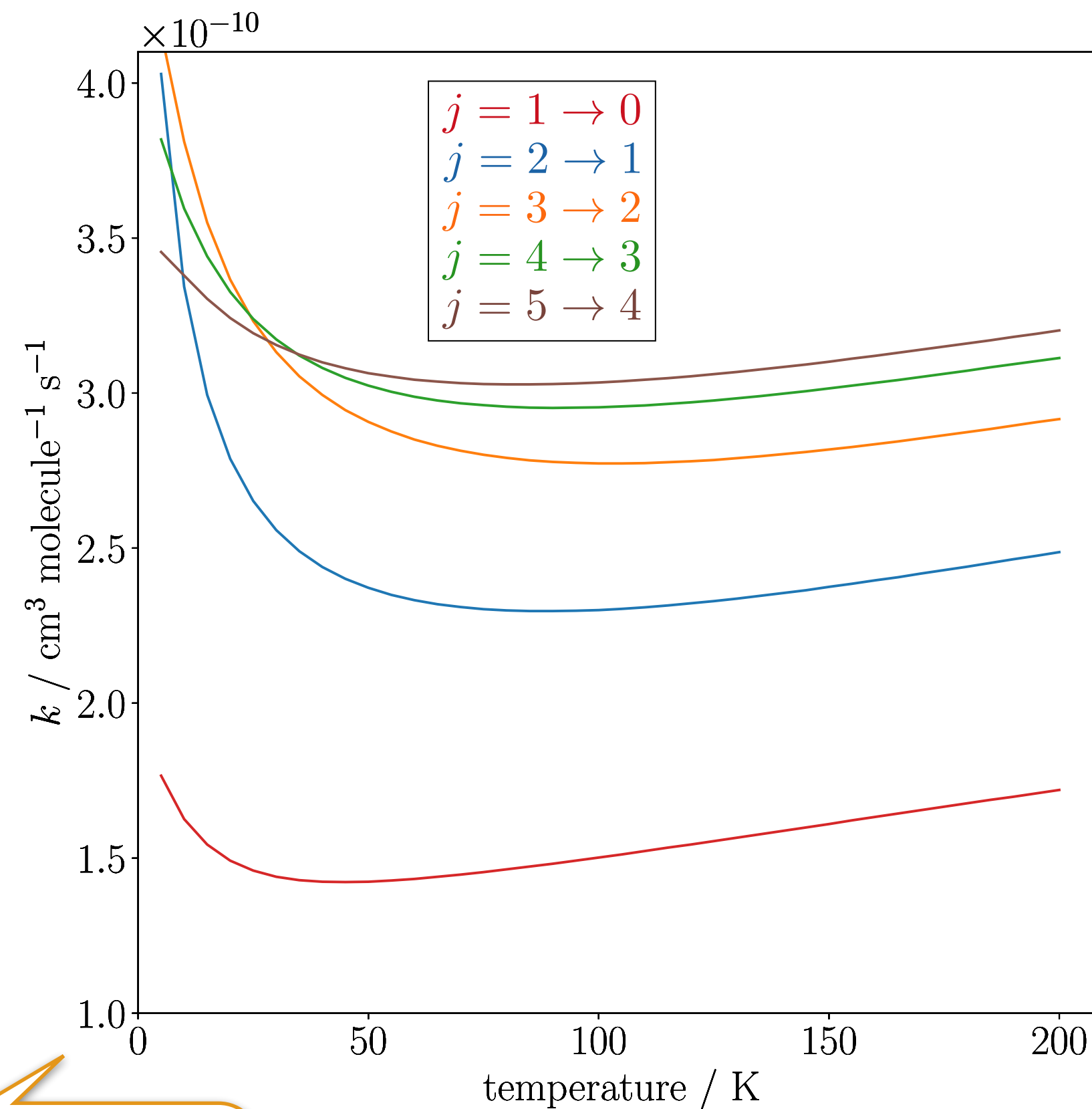
$$\sigma_{j_B \rightarrow j'_B}(E_c) = \frac{\pi}{k_{j_B}^2} \sum_{p,J,M} \frac{2J+1}{2j_B+1} \sum_{l,l'} \left| \delta_{ll'} - S_{ll'}^{pJM} \right|^2 \longrightarrow k_{j_B \rightarrow j'_B}(T) = \left(\frac{8}{\pi \mu k_B^3 T^3} \right)^{1/2} \int_0^\infty \sigma_{j_B \rightarrow j'_B}(E_c) E_c \exp(-E_c/k_B T) dE_c$$



Tonolo+, *MNRAS*, **527** (2024).

close coupling calculations

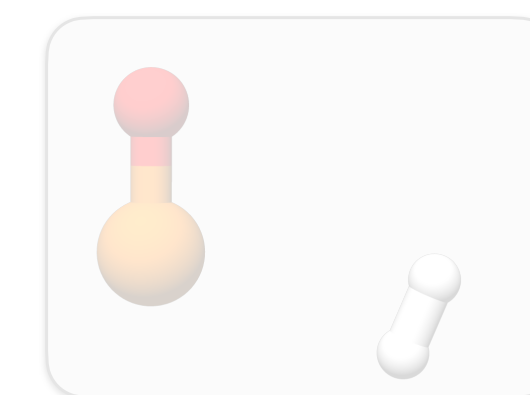
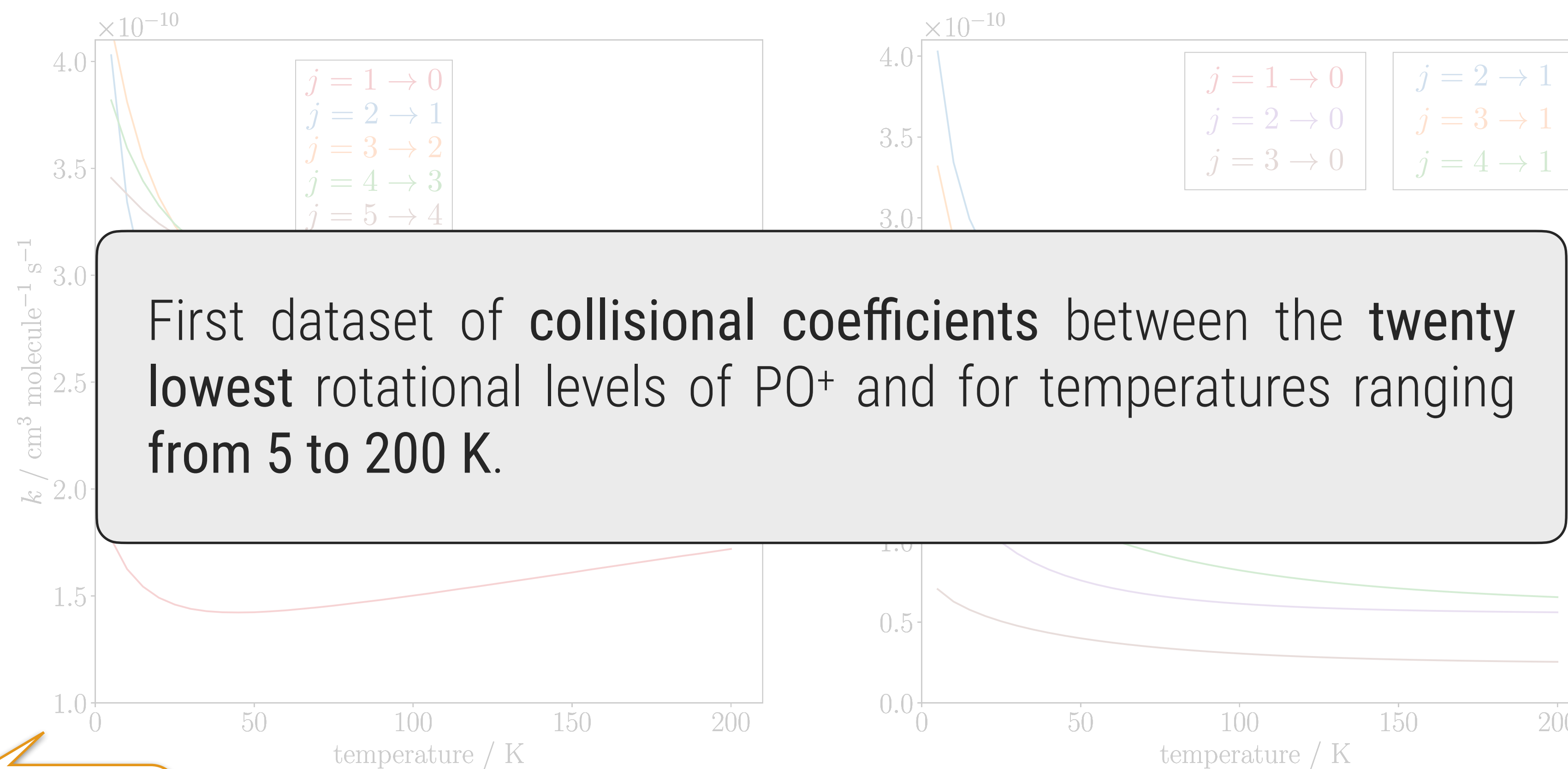
$$\sigma_{j_B \rightarrow j'_B}(E_c) = \frac{\pi}{k_{j_B}^2} \sum_{p,J,M} \frac{2J+1}{2j_B+1} \sum_{l,l'} \left| \delta_{ll'} - S_{ll'}^{pJM} \right|^2 \longrightarrow k_{j_B \rightarrow j'_B}(T) = \left(\frac{8}{\pi \mu k_B^3 T^3} \right)^{1/2} \int_0^\infty \sigma_{j_B \rightarrow j'_B}(E_c) E_c \exp(-E_c/k_B T) dE_c$$



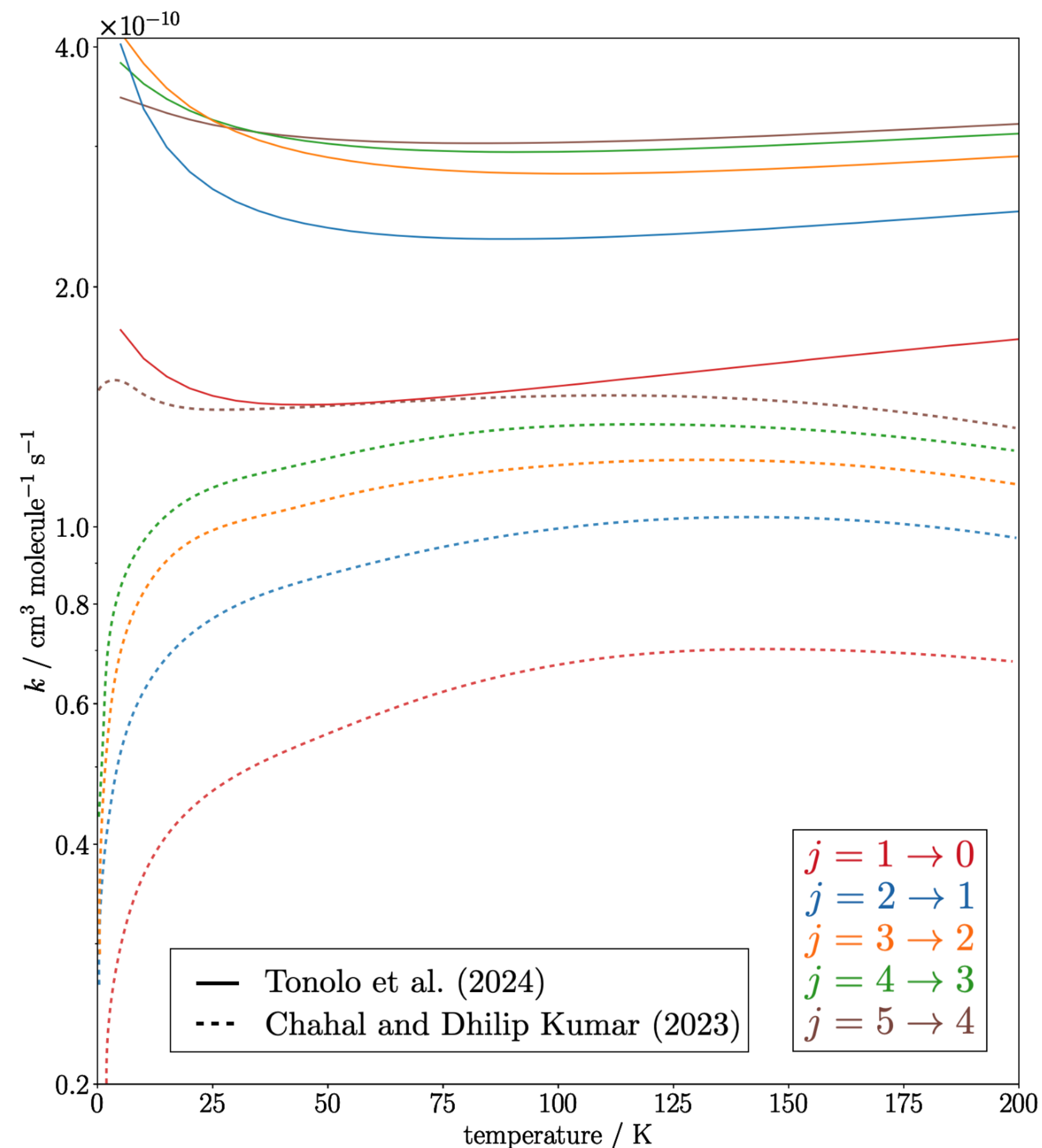
Tonolo+, *MNRAS*, **527** (2024).

close coupling calculations

$$\sigma_{j_B \rightarrow j'_B}(E_c) = \frac{\pi}{k_{j_B}^2} \sum_{p,J,M} \frac{2J+1}{2j_B+1} \sum_{l,l'} \left| \delta_{ll'} - S_{ll'}^{pJM} \right|^2 \longrightarrow k_{j_B \rightarrow j'_B}(T) = \left(\frac{8}{\pi \mu k_B^3 T^3} \right)^{1/2} \int_0^\infty \sigma_{j_B \rightarrow j'_B}(E_c) E_c \exp(-E_c/k_B T) dE_c$$



Why not using the He collider?



Comparison of some rotational de-excitation rate coefficients: **mass-scaled** results from previous calculations on the PO⁺ and He system (dashed lines) vs. this work (solid lines).



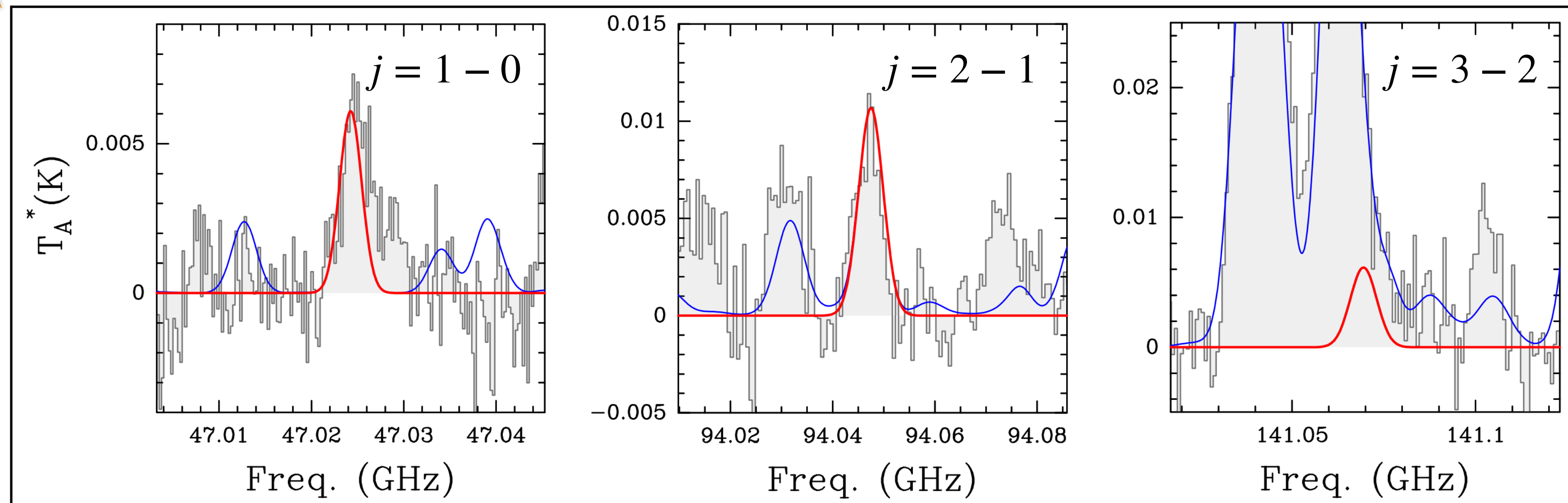
Inaccuracy of He as a proxy of the H₂ perturber!

Tonolo+, *MNRAS*, **527** (2024);
Chahal+, *MNRAS*, **523** (2023).

Radiative transfer modeling

The recent detection of PO⁺ in the **molecular cloud G+0.693-0.02** where, due to the low density of H₂, **LTE condition are not achieved**, motivated the need of collisional rate coefficients to properly model its abundance.

V. M. Rivilla, et al. (2022), *FSPAS*, **9**, 829288.



G+0.693-0.027 molecular cloud:

- Very low densities of H₂ ($\sim 1 \times 10^4 \text{ cm}^{-3}$);
- High kinetic temperatures ($\sim 150 \text{ K}$);

Quasi-thermalization condition!

$$T_{\text{ex}} \sim 5 - 20 \text{ K}$$

S. Zeng, et al. (2018), *MNRAS*, **478**, 2962;

S. Zeng, et al. (2020), *MNRAS*, **497**, 4896.

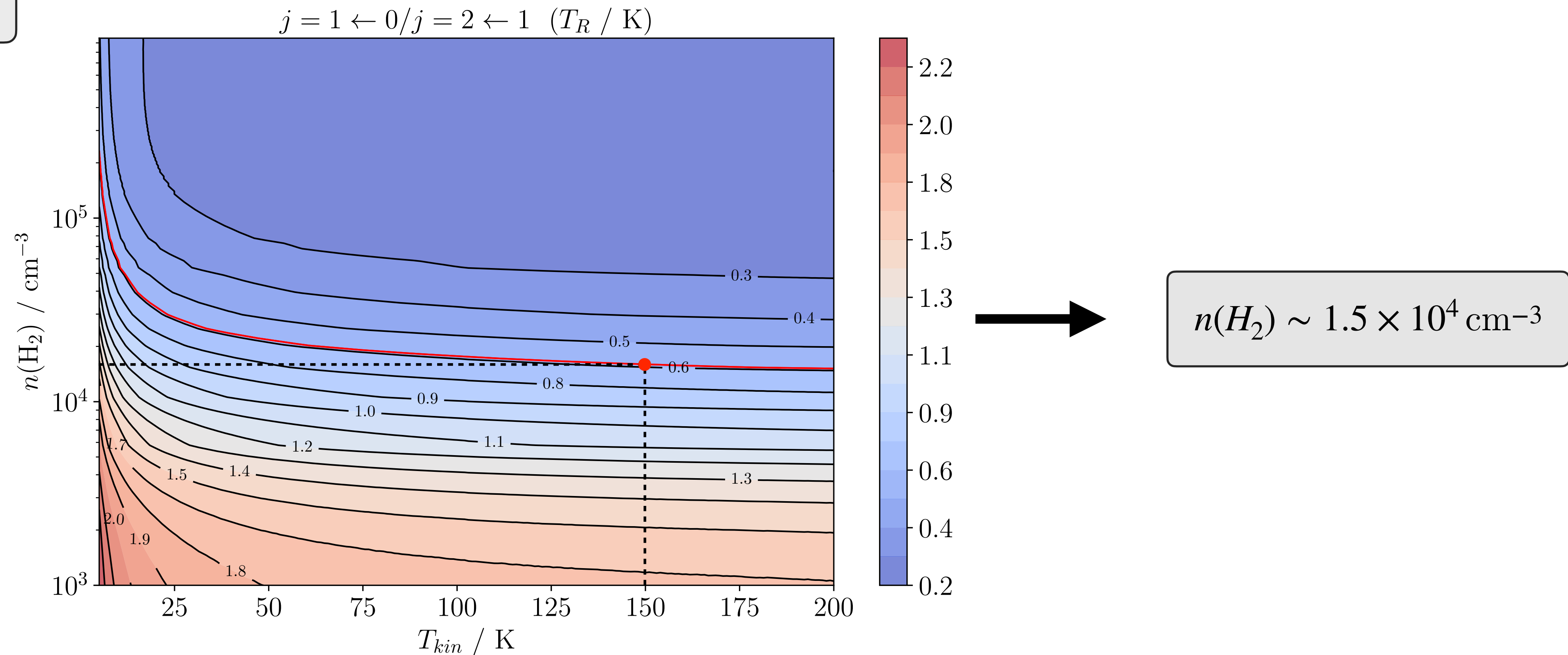
P. F. Goldsmith and W. D. Langer (1999), *ApJ*, **517**, 209.

Radiative transfer modeling



RADEX

Radiative transfer calculations have been used to model the observation of PO⁺ in G+0.693-0.027.



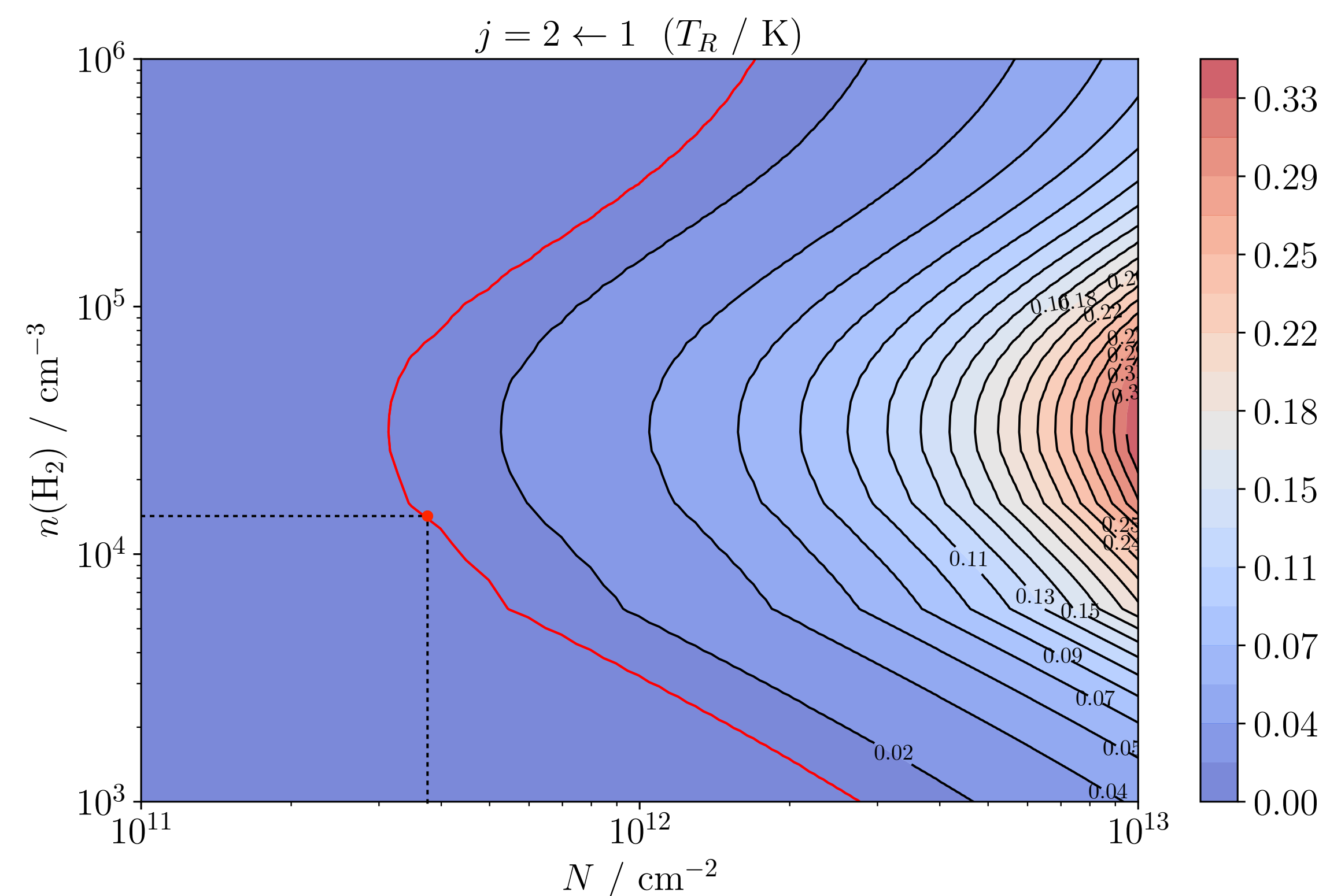
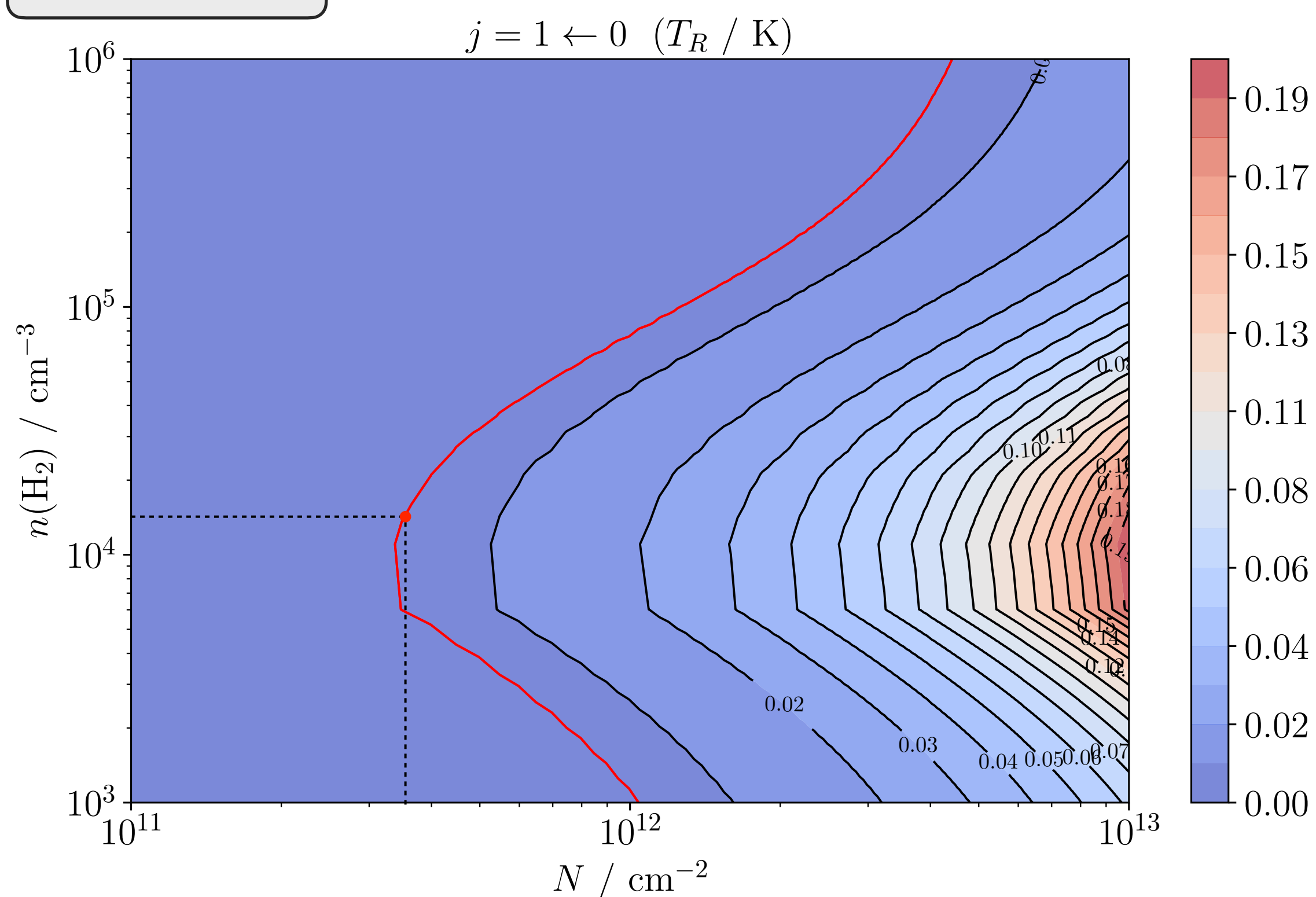
Variation of the $j=1-0/j=2-1$ intensity ratio as a function of T_{kin} and $n(H_2)$. The observed intensity ratio is highlighted in red.

Radiative transfer modeling

Single transition analysis: intensity trends at 150 K as a function of $n(\text{H}_2)$ and N .



RADEX



$N \sim 3.7 \times 10^{11} \text{ cm}^{-2}$



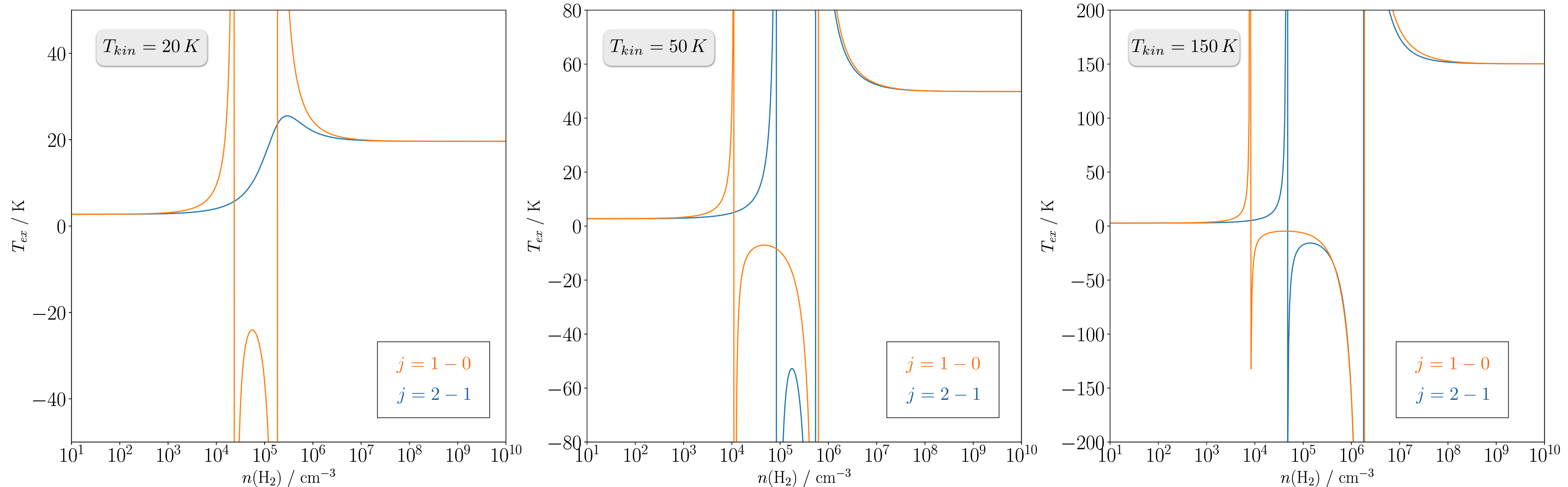
A factor of $\sim 60\%$ smaller than the previously LTE-derived value.

Radiative transfer modeling



Behavior of T_{ex} for $j=1-0$ and $j=2-1$ at different densities and temperature conditions → For densities around 10^5 cm^{-3} , **maser phenomena** are predicted.

RADEX

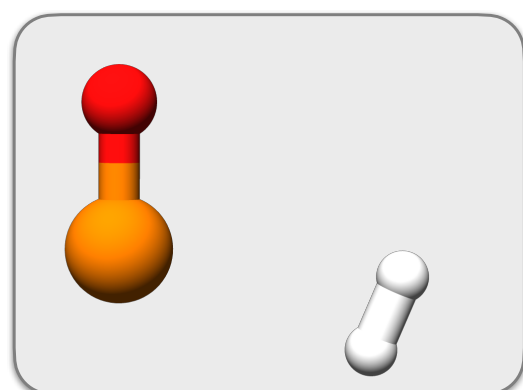


Variation of the excitation temperature as a function of $n(\text{H}_2)$ for the first two transitions of PO⁺ at $T_{kin}=20, 50, 150 \text{ K}$.

Tonolo+, *MNRAS*, **527** (2024).

Conclusions

PO⁺ and H₂
collisional system



Characterization of the Interaction Potential

Single point energies at CCSD(T)-F12/aug-cc-pV(Q+d)Z level of theory + **spherical average approximation** to fit the potential.

Scattering Calculations

First dataset of **collisional coefficients** between the **twenty lowest** rotational levels of PO⁺ and for temperatures ranging from 5 to 200 K.

Astrophysical Modeling

Radiative transfer modeling of the recently observed lines that allowed to **refine the abundance of PO⁺** measured in the G+0.693–0.027 cloud. Strong **maser effects** highlight the importance of collisional data for astrophysical applications.

Thank you!



Prof. François Lique



Prof. Cristina Puzzarini
Dr. Luca Bizzocchi
Dr. Mattia Melosso



Dr. Victor M. Rivilla

